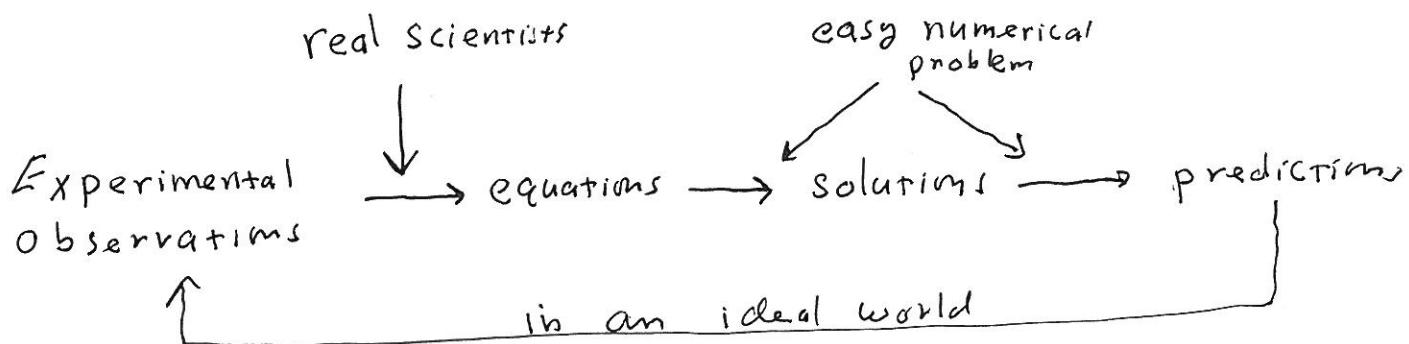


(1)

10/19/04



Two things changed this view:

1. The open question is: how do low level eqns. (e.g. Ht) describe high level behavior (e.g. people).
2. "Easy" numerical problem isn't. Often easy should be replaced by impossible.

So, people started studying equations and their solutions. (Dynamical systems a.k.a. chaos theory).

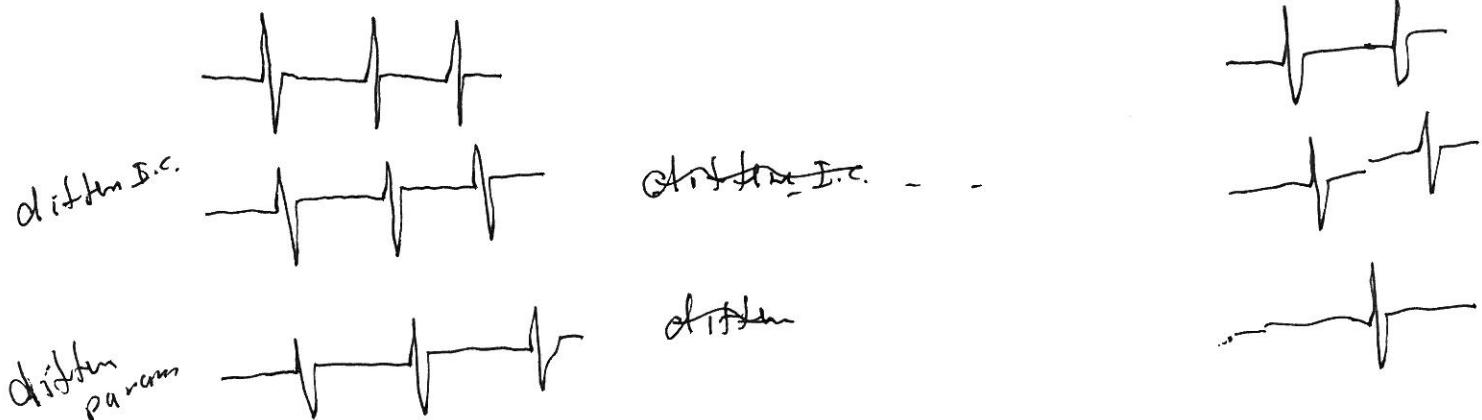
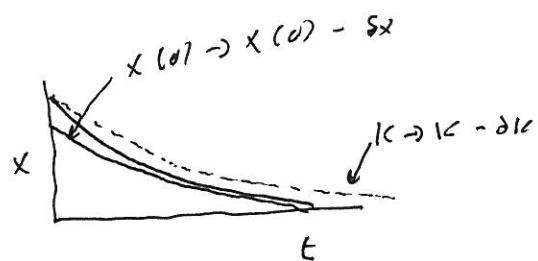
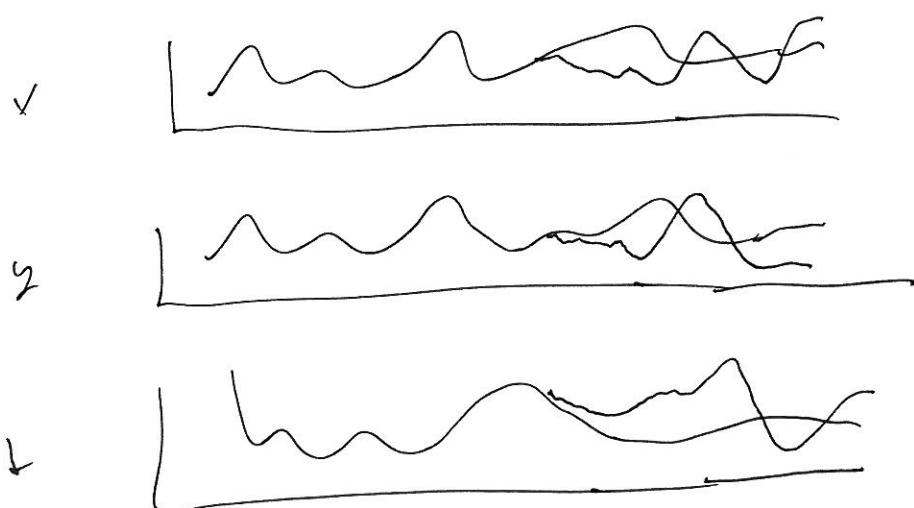
Questions:

1. Do equations fall into distinct classes, w/ each class having characteristic solutions?
Somewhat
2. What equation can be solved?
depends on hand
3. What does question (2) mean?
hand
4. Are solutions to equations sensitive to
 - a) initial condition?
sensitivity
 - b) parameters? or eqn
almost always

(2)

Example:

$$\frac{dx}{dt} = -kx$$

Edward Lorenz (~ 1960)

~~cha~~

Birth of a
freak

(3)

We're going to use some of these techniques to examine spiking neurons, starting with HH.

$$T \frac{dV}{dt} = F(V, m, h, n) + I$$

$$F(V, m, h, n) = -(\bar{g}_L(V - \Sigma_L)) + \bar{g}_{Na} m^3 h (V - \Sigma_{Na}) - \bar{g}_{K} n^4 (V - \Sigma_K)$$

$$T = \frac{C_m}{\bar{g}_L} = C_m R_m$$

$$\bar{g}_{Na} = \frac{\bar{g}_{Na}}{\bar{g}_L} \quad \bar{g}_K = \frac{\bar{g}_K}{\bar{g}_L}$$

$$T_m(V) \frac{dm}{dt} = M_\infty(V) - m$$

$$T_h(V) \frac{dh}{dt} = H_\infty(V) - h$$

$$T_n(V) \frac{dn}{dt} = N_\infty(V) - n$$

H-P or E!!

④

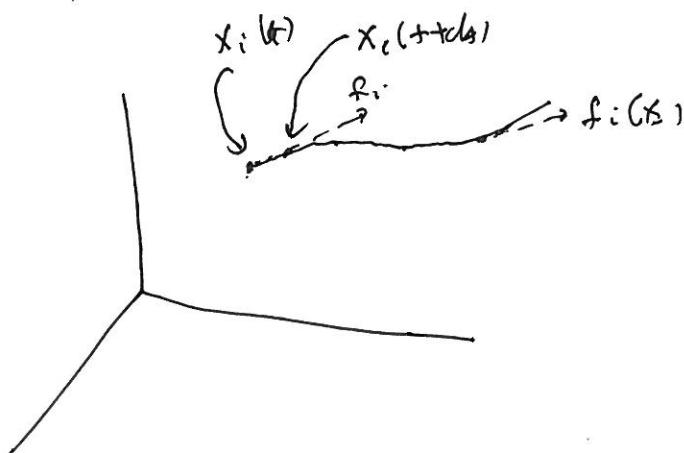
Fact

Aside on solutions to equations

$$\frac{dx_i}{dt} = f_i(x) \quad \underbrace{x_1, x_2, \dots}$$

$$x_i(t+dt) \approx x_i(t) + \frac{dx_i}{dt} dt$$

$$= x_i(t) + f_i(x) dt$$



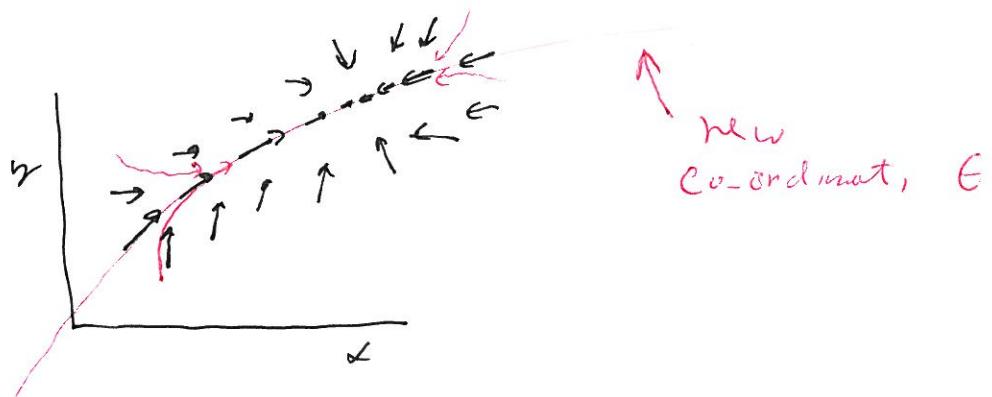
- Solutions are trajectories
- f is vector field

Problem: go from vector field to trajectories

Fact

(5)

Dimensionality reduction



$$\frac{dx}{dt} = f_x(x, y)$$

$$\frac{dy}{dt} = f_y(x, y) \Rightarrow \frac{d\theta}{dt} = F(\theta)$$

The plan: find strongly attracting attractors

∞ -dimensional ODE's (aka PDE's) can have finite dimensional attractors!

- reduce 4-D HH \rightarrow 2-D

(6)

Approximation # 1: $\gamma \gg T_m, T_h, T_n$ (high resistance cell).

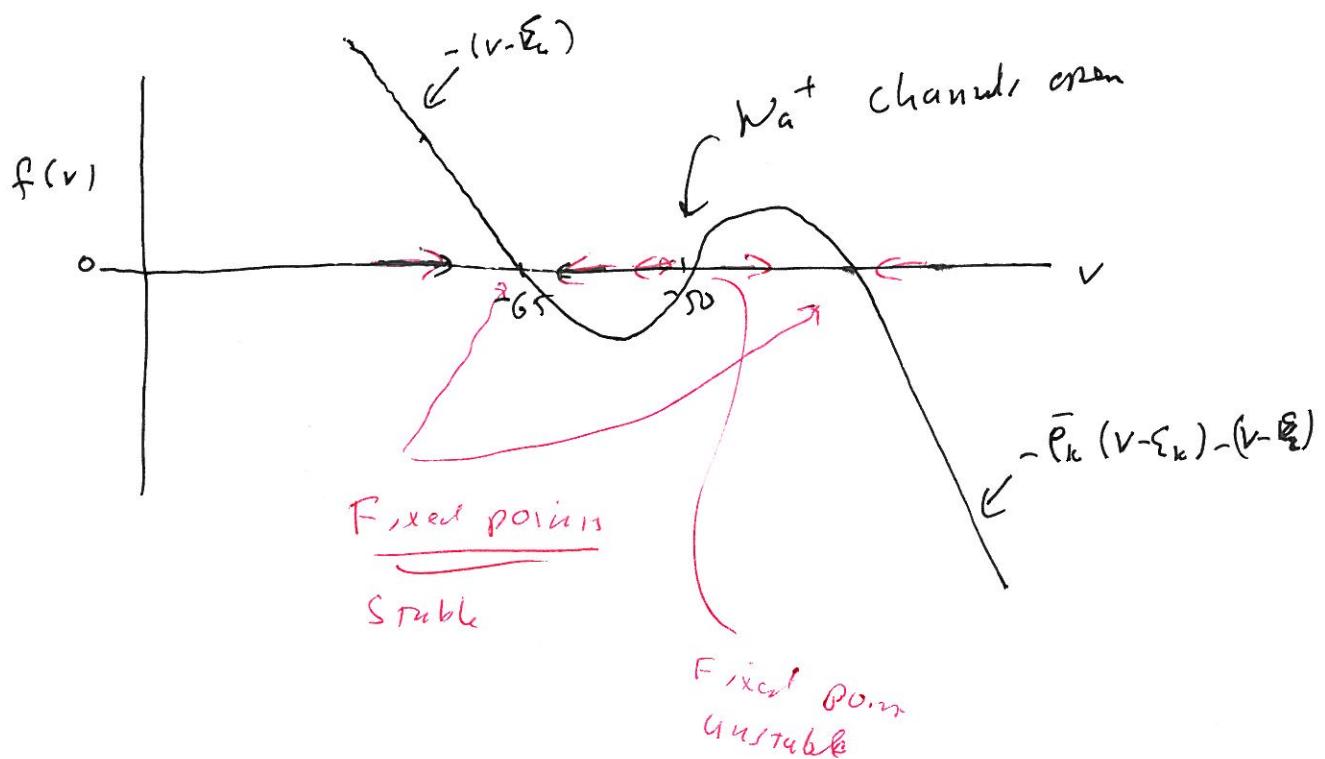
$$\Rightarrow m \approx m_\infty(v)$$

$$n \approx n_\infty(v)$$

$$h \approx h_\infty(v)$$

$$\begin{aligned} \gamma \frac{dv}{dt} &= F(v, m_\infty(v), h_\infty(v), n_\infty(v)) + I \\ &= f(v) + I \end{aligned}$$

What does $f(v)$ look like?



(7)

$$\nabla \frac{dv}{ds} = f(v) + I$$

$$f(v_0) + I = 0 \quad (\text{fixed point})$$

$$v = v_0 + \delta v, \quad \delta v \text{ small}$$

$$\nabla \frac{d}{ds} V = \nabla \frac{d}{ds} (V_0 + \delta v) = \nabla \frac{d \delta v}{ds} = f(v_0 + \delta v) + I \\ = f(v_0) + \underbrace{\delta v f'(v_0) + I}_{f'(v_0) \infty}$$

$$\nabla \frac{d \delta v}{ds} = f'(v_0) \delta v \Rightarrow \boxed{\delta v(+) = \delta v(0) e^{f'(v_0)s}}$$

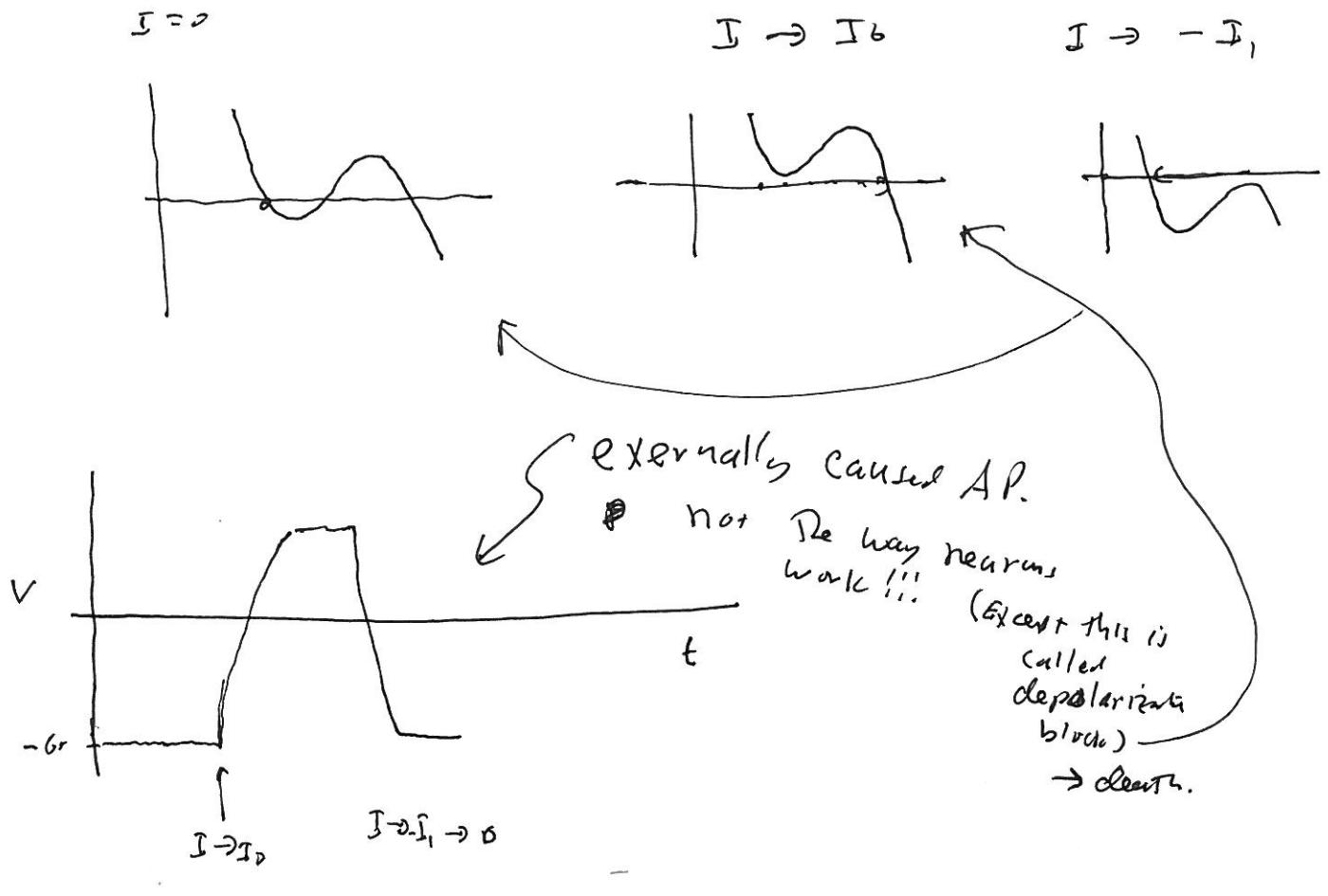
$$f'(v_0) < 0 \Rightarrow \text{stable}$$

$$f'(v_0) > 0 \Rightarrow \text{unstable}$$

III

(8)

No repetitive firings !!!



Approx. #1 was a good idea, but it didn't work!!

- that "explains" why $T \sim T_m, T_h, T_n \dots$

(9)

Approx. # 2:

$$m = m_\infty(v)$$

$$h = h_\infty(u)$$

$$n = n_\infty(u')$$

$$\frac{dh}{dt} = \frac{1}{T_h} (h_\infty(v) - h) = \frac{dh_\infty(u)}{du} \frac{du}{dt}$$

$$\frac{du}{dt} = \frac{\frac{h_\infty(v) - h(u)}{T_h(v)}}{\frac{dh_\infty(u)}{du}} \quad \begin{matrix} v = u + V - u \\ \longrightarrow \end{matrix} \quad \frac{V-u}{T_h(v)}$$

$$\frac{du'}{dt} = \frac{\frac{n_\infty(v) - n(u')}{T_n(v)}}{\frac{dn_\infty(u)}{du}} \quad \longrightarrow \quad \frac{V-u'}{T_n(v)}$$

u, u' relaxes toward V w/ open time constants
 $T_h(v), T_n(v)$.

$$\frac{d}{dt}(u - u') = \frac{V}{T_h} - \frac{V}{T_n} - \left(\frac{u}{T_h} - \frac{u'}{T_n} \right)$$

one app

(10)

$$\gamma \frac{dv}{dt} = F(v, m_\infty(v), h_\infty(u), n_\infty(u)) + I$$

$F_{\text{true}}(v, m_\infty(v), h, n) \leftarrow \text{true}$

~~Kleiner Abboott & Kepler:~~

$$\frac{dF_{\text{true}}}{dt} \Big|_{V \neq v_i, n \neq n_i} = \frac{dF}{dt} \Big|_{V \neq v_i, n \neq n_i}$$

$n = h_\infty,$
 $v = v_\infty$

$$= \frac{\partial F}{\partial h} \frac{dh}{dt} + \frac{\partial F}{\partial n} \frac{dn}{dt} = \left[\frac{\partial F}{\partial h_\infty} \frac{dh_\infty}{dt} + \frac{\partial F}{\partial n_\infty} \frac{dn_\infty}{dt} \right] \frac{dy}{dt}$$

$$\frac{dy}{dt} = \frac{\frac{\partial F}{\partial h_\infty(u)} \frac{h_\infty(v) - h_\infty(u)}{T_h(v)} + \frac{\partial F}{\partial n_\infty(u)} \frac{n_\infty(v) - n_\infty(u)}{T_h(v)}}{\frac{\partial F}{\partial h_\infty} \frac{dh_\infty}{dt} + \frac{\partial F}{\partial n_\infty} \frac{dn_\infty}{dt}}$$

$$\frac{du}{dt} = g(v, u)$$

(11)

$$T \frac{dv}{dt} = f(v, u) + I$$

$$\frac{du}{dt} = g(v, u)$$

$g(0, 0) = 0$ $v=u$ is equilibrium.

~~~~~

### nullcline analysis:

1<sup>st</sup> question: what are the fixed points???

$$\frac{dv}{dt} = \frac{du}{dt} = 0$$

nullclines:

$$\left. \begin{array}{l} \frac{dv}{dt} = 0 \\ \frac{du}{dt} = 0 \end{array} \right\}$$

define curves  
in 2-D space;  
intersection are fixed  
points.

(12)

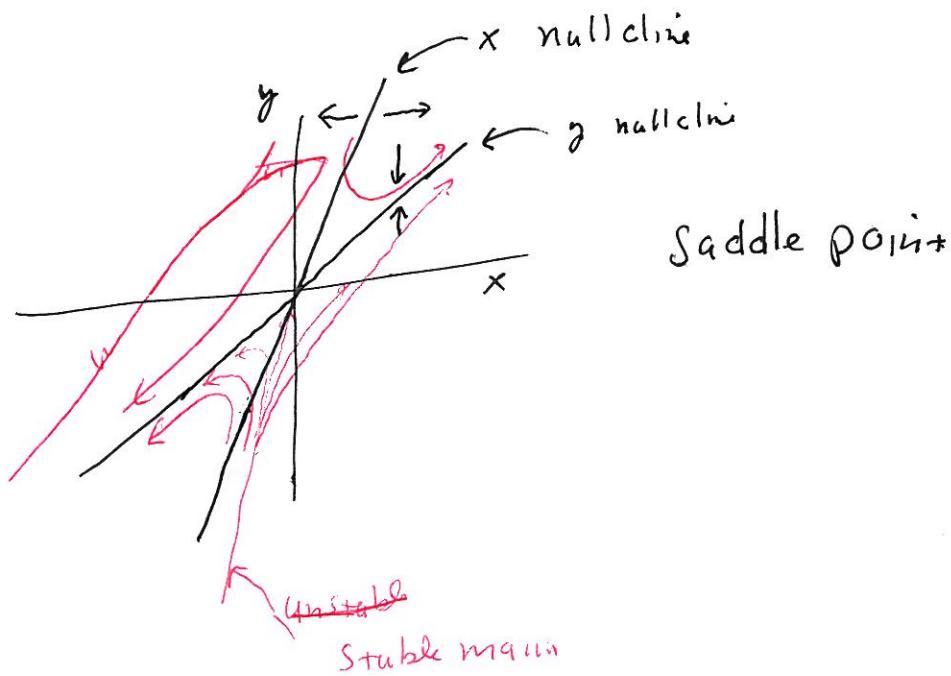
Example

$$\frac{dx}{dt} = ax - by$$

$$\frac{dy}{dt} = cx - dy$$

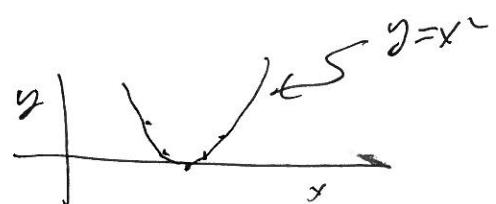
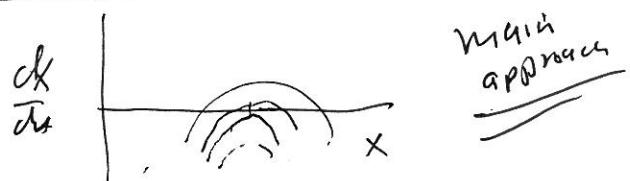
$$\frac{dx}{dt} = 0 \Rightarrow y = \frac{a}{b}x$$

$$\frac{dy}{dt} = 0 \Rightarrow y = \frac{c}{d}x$$

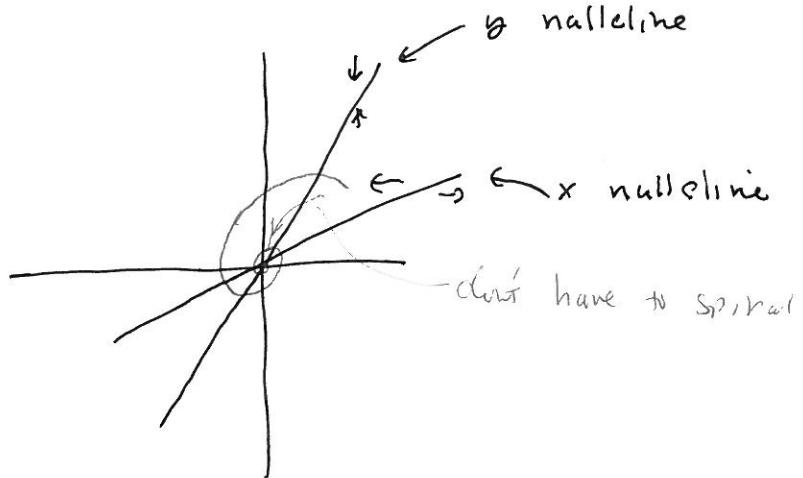


$$\frac{dx}{dt} = y - x^2$$

$$\frac{dy}{dt} = x - y$$



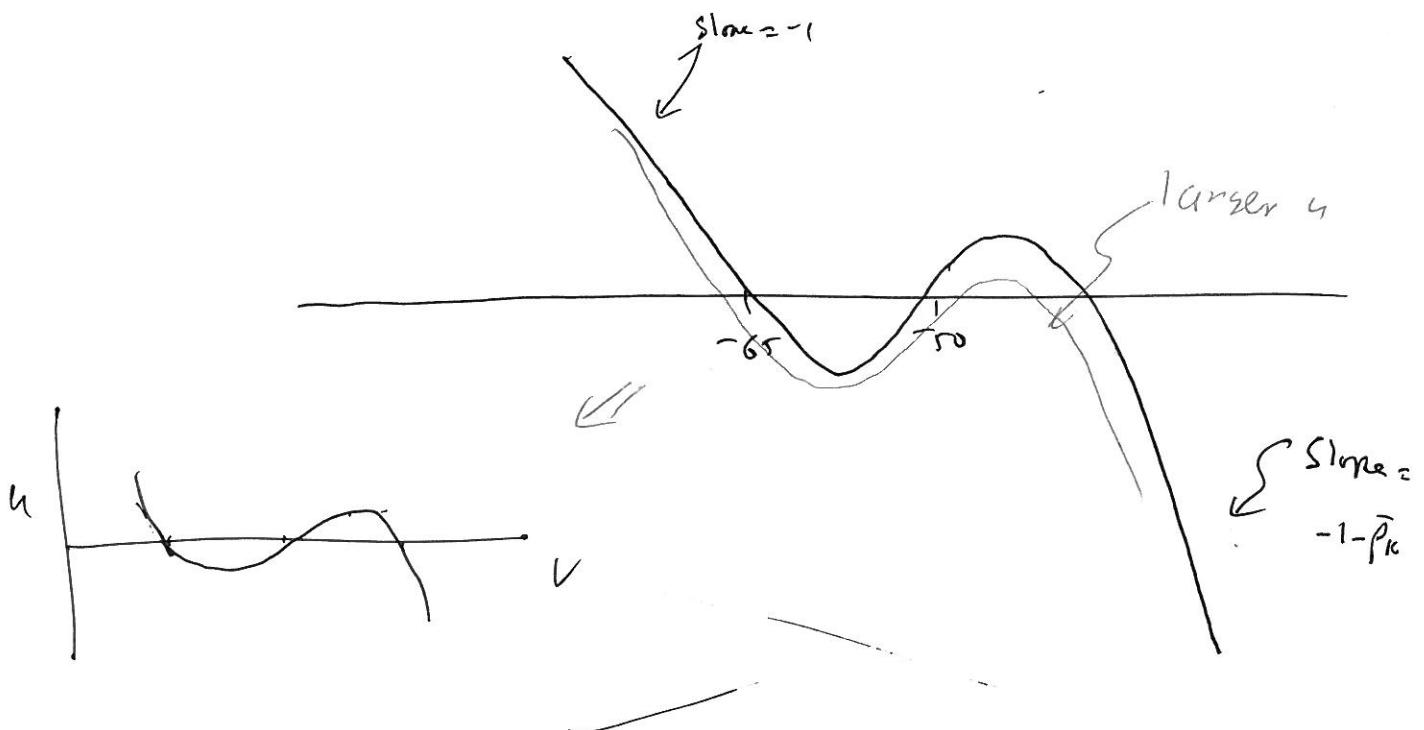
(13)



|||||

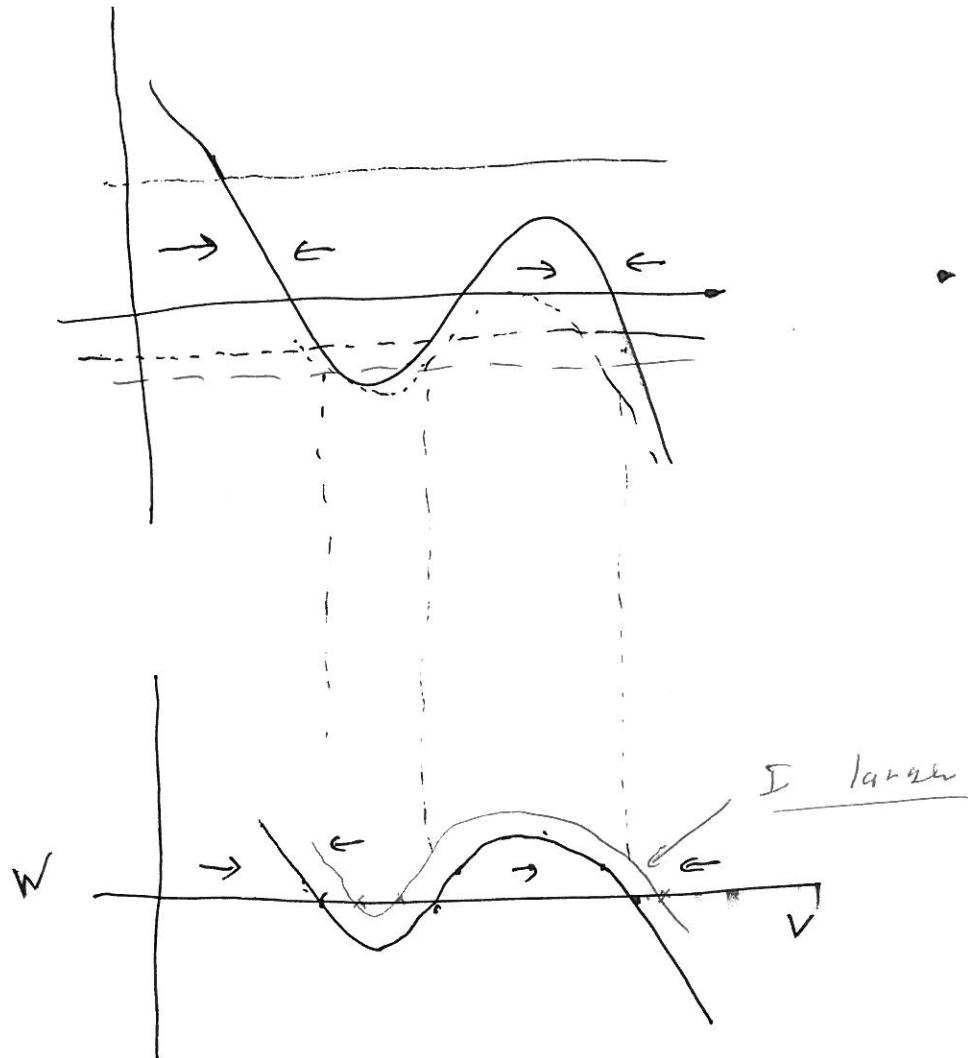
What does  $f(v, u)$  look like?

$$f(v, u) = -(v - \Sigma_L) - \bar{\rho}_{w_a} m_\infty^3(v) h_\infty(u) (v - \Sigma_{w_a}) - \bar{\rho}_{l_c} n_\infty^4(u) (v - \Sigma_{l_c})$$



(14)

$$\frac{dv}{dt} = 0 \Rightarrow f(v, w) \stackrel{?}{=} I = 0$$



(15)

FitzHugh - Nagumo:

$$\dot{V} = V - \frac{V^3}{3} - W + I$$

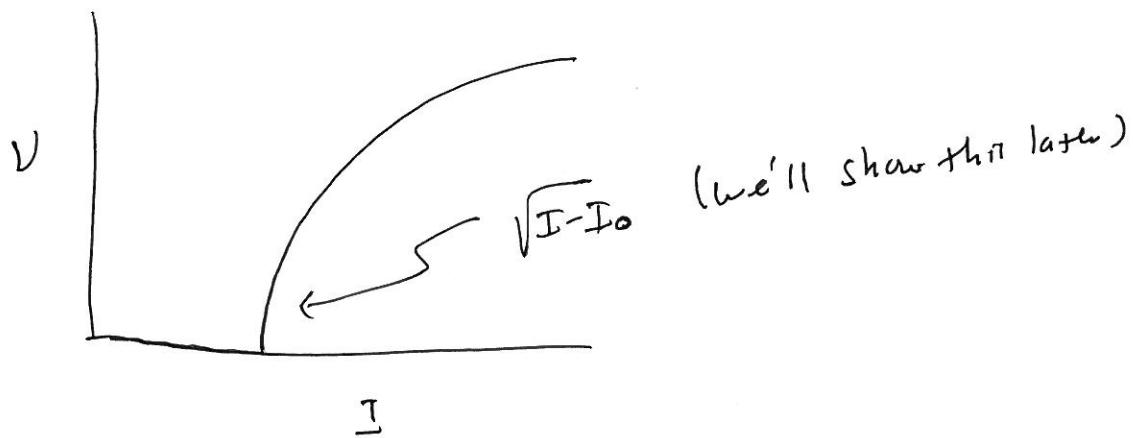
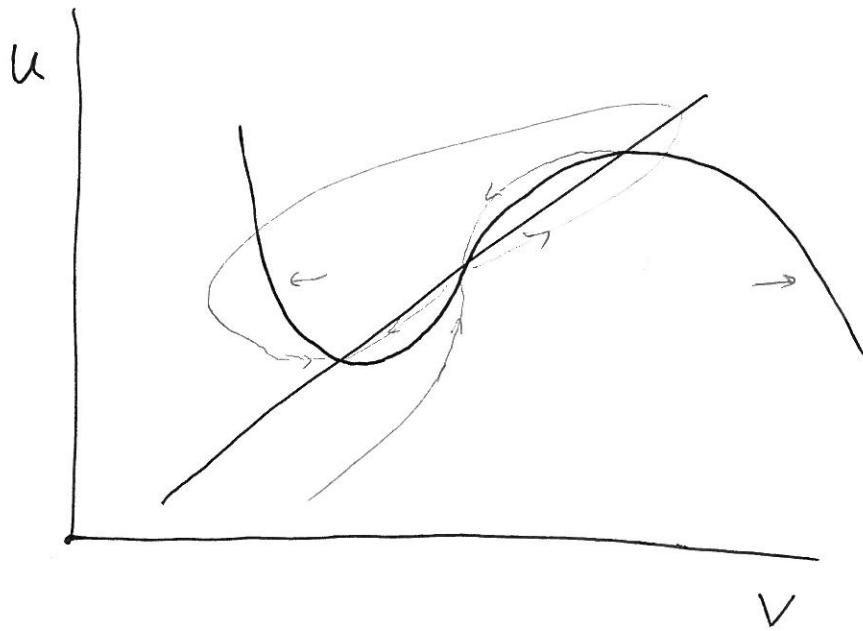
$$\dot{W} = \phi(V + a - bW)$$

Morris - Lecar

$$C_m \frac{dV}{dt} = - \underbrace{\sum_{Ca} \bar{g}_{Ca} M_\infty(V) (V - \Sigma_{Ca})}_{-\bar{g}_L (V - \Sigma_L)} - \bar{g}_K W (V - \Sigma_K) - \bar{g}_I (V - \Sigma_I) + I$$

$$T_w(V) \frac{dW}{dt} = \frac{W \bar{g}(V) - w}{T_w(V)}$$

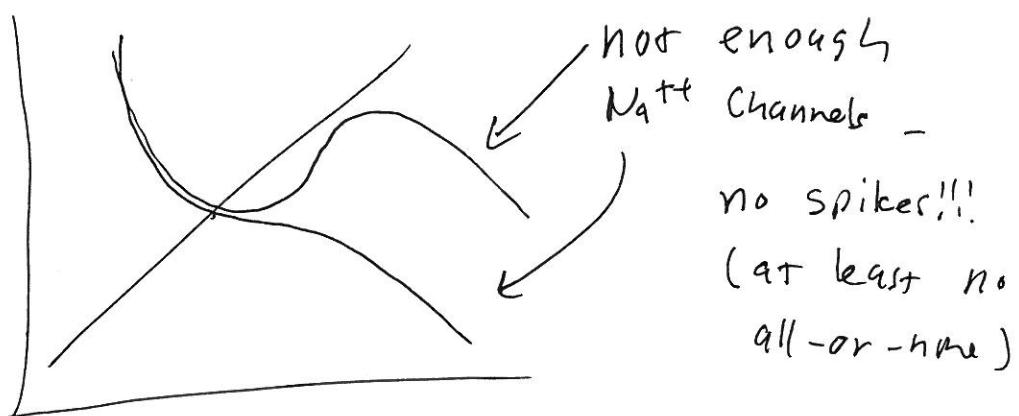
(15)

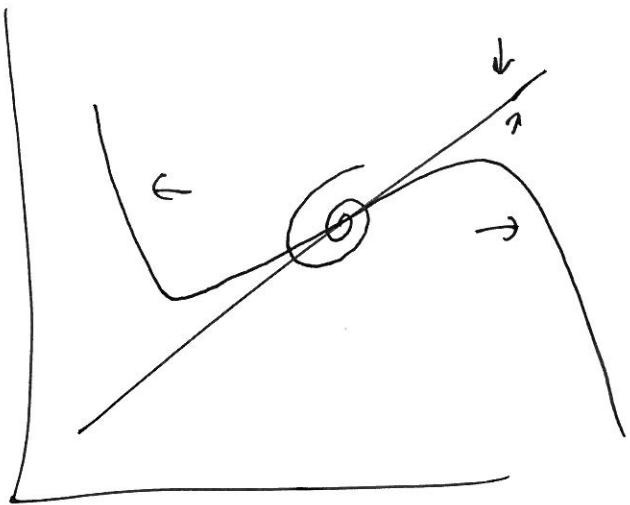


Type II!!

all-or-none

(17)





Hopf - eigenvalues - subcritical

